

Hamiltonian, and operator-definition of angular momentum for central potential,

$$H = \frac{p^2}{2m} + V(r) = -\frac{\hbar^2}{2m} \nabla^2 + V(r) \quad \mathbf{L} = \mathbf{r} \times \mathbf{p} = \mathbf{r} \times \frac{\hbar}{i} \nabla \quad [\text{I.1}]$$

Some commutators that are easy to see given the canonical $[r_i, p_j] \equiv i\hbar \cdot \delta_{ij}$,

$$[L_x, y] = [yp_z - zp_y, y] = -z[p_y, y] = i\hbar z \quad [L_x, p_y] = [yp_z - zp_y, p_y] = [y, p_y]p_z = i\hbar p_z \quad [L_z, x] = 0, \quad [L_x, p_x] = 0 \quad [\text{I.2}]$$

$$[L_x, L_y] = [L_x, zp_x - xp_z] = [L_x, z]p_x - x[L_x, p_z] = -i\hbar yp_x + i\hbar xp_y = i\hbar L_z \quad [\text{I.3}]$$

By cyclic permutation of [I.3], we can deduce,

$$[L_i, L_j] \equiv i\hbar L_k \quad \leftrightarrow \quad \mathbf{L} \times \mathbf{L} = i\hbar \cdot \mathbf{L} \quad [\text{I.4}]$$

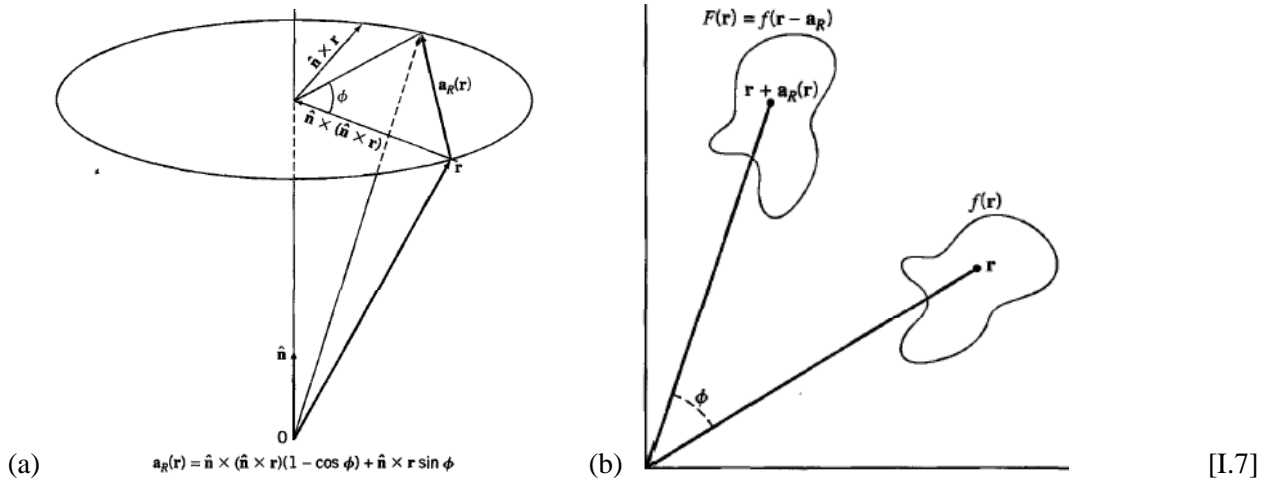
Analogy with rigid rotations

Consider f to be a function of a vector. It's mathematics, not quantum mechanics. Then, finite and infinitesimal translations are effected by linearity (Taylor series),

$$f(\mathbf{r}) \mapsto F(\mathbf{r} + \mathbf{a}) = f(\mathbf{r}) \quad F(\mathbf{r} + \mathbf{\epsilon}) = F(\mathbf{r}) + \mathbf{\epsilon} \cdot \nabla F(\mathbf{r}) = f(\mathbf{r}) \quad [\text{I.5}]$$

A finite *rotation*, on the other hand, is a bit more complicated. Recall QM 08, 320 - pr 4-3 - generator of arbitrary rotations in R3 and/or CM 04, 181 - de 12 - the rotation formula, where we derived,

$$\mathbf{r} \rightarrow \mathbf{r}' = \mathbf{n} \times (\mathbf{n} \times \mathbf{r})(1 - \cos \phi) - \mathbf{n} \times \mathbf{r} \cdot \sin \phi = \mathbf{a}_R \quad [\text{I.6}]$$



Rotation about an axis defined by the unit vector $\hat{\mathbf{n}}$ and the rotation angle ϕ (a) shows the displacement \mathbf{a}_R (see [I.6]) of the point whose position vector is \mathbf{r} . (b) illustrates the active rotation of a function or state $f(\mathbf{r})$ about an axis ($\hat{\mathbf{n}}$) perpendicular to the plane of the figure: $f(\mathbf{r}) \rightarrow F(\mathbf{r}) = f(\mathbf{r} - \mathbf{a}_R)$.

Verify $\mathbf{r} \rightarrow \mathbf{r}' = \mathbf{n} \times (\mathbf{n} \times \mathbf{r})(1 - \cos \phi) - \mathbf{n} \times \mathbf{r} \cdot \sin \phi = \mathbf{a}_R$.

See QM 08, 320 - pr 4-3 - generator of arbitrary rotations in R3 and/or CM 04, 181 - de 12 - the rotation formula.

Show that it gives the expected answer for a rotation about the z -axis. Using $\hat{\mathbf{z}} \times (\hat{\mathbf{z}} \times (x, y)) = \hat{\mathbf{z}} \times (+y, -x) = (+y, -x)$, an arbitrary vector (x, y) in the xy -plane is rotated as,

$$\begin{aligned}
 (x, y) \rightarrow (x', y') &= \hat{z} \times (\hat{z} \times (x, y))(1 - \cos \phi) - \hat{z} \times (x, y) \cdot \sin \phi = \hat{z} \times (+y, -x)(1 - \cos \phi) - (+y, -x) \cdot \sin \phi \\
 &= \begin{bmatrix} -x \\ -y \end{bmatrix} (1 - \cos \phi) - \begin{bmatrix} y \\ -x \end{bmatrix} \sin \phi = \begin{bmatrix} x \\ y \end{bmatrix} \cos \phi + \begin{bmatrix} -y \\ x \end{bmatrix} \sin \phi - \begin{bmatrix} x \\ y \end{bmatrix} = \boxed{\begin{bmatrix} \cos \phi & -\sin \phi \\ \sin \phi & \cos \phi \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} - \begin{bmatrix} x \\ y \end{bmatrix}} \quad [\text{I.1}]
 \end{aligned}$$

Mistake? I think I'm off by a minus sign in [I.1]. I think the $-\begin{bmatrix} x \\ y \end{bmatrix}$ should be a $+\begin{bmatrix} x \\ y \end{bmatrix}$ (the very last column-vector that looks like a translation added onto the rotation we derived, in [I.1]. If this was a $+$, it would look a *bit* like the end I'm working towards in Shankar 12.4.4, with the cosines and sines constituting broken-up components of $\vec{\theta}$ that is supposed to appear in a cross product).