

Nice proof here. The rigorous proof, actually.

For a plane movement

$$\vec{r}(t) = x(t)\vec{i} + y(t)\vec{j}$$

$$\vec{v}(t) = x'(t)\vec{i} + y'(t)\vec{j}$$

$$\vec{a}(t) = x''(t)\vec{i} + y''(t)\vec{j}$$

With the modulus

$$|\vec{r}(t)| = +\sqrt{x^2(t) + y^2(t)}$$

$$|\vec{v}(t)| = +\sqrt{(x'(t))^2 + (y'(t))^2}$$

$$|\vec{a}(t)| = +\sqrt{(x''(t))^2 + (y''(t))^2}$$

Okay?

Now assume that u want to express the same vector (i.e. the acceleration) in polar plane coordinates ρ and φ , but in the same (\vec{i}, \vec{j}) basis.

U know that

$$x(t) = \rho(t) \cos \varphi(t)$$

$$y(t) = \rho(t) \sin \varphi(t)$$

U'll simply compute that

$$x'(t) = (\rho(t) \cos \varphi(t))' = \rho'(t) \cos \varphi(t) - \rho(t) \varphi'(t) \sin \varphi(t)$$

$$y'(t) = (\rho(t) \sin \varphi(t))' = \rho'(t) \sin \varphi(t) + \rho(t) \varphi'(t) \cos \varphi(t)$$

and

$$\begin{aligned} x''(t) &= (x'(t))' = \rho''(t) \cos \varphi(t) - 2\rho'(t) \varphi'(t) \sin \varphi(t) - \rho(t) \varphi''(t) \sin \varphi(t) - \\ &\quad \rho(t) (\varphi'(t))^2 \cos \varphi(t) \\ y''(t) &= (y'(t))' = \rho''(t) \sin \varphi(t) + 2\rho'(t) \varphi'(t) \cos \varphi(t) + \rho(t) \varphi''(t) \cos \varphi(t) - \\ &\quad \rho(t) (\varphi'(t))^2 \sin \varphi(t) \end{aligned}$$

Assume a circular movement. That means $\rho(t) = \text{const.} = R$, where R is the radius' circle. Then

$$x'(t) = -R\varphi'(t) \sin \varphi(t)$$

$$y'(t) = R\varphi'(t) \cos \varphi(t)$$

$$x''(t) = -R\varphi''(t) \sin \varphi(t) - R(\varphi'(t))^2 \cos \varphi(t)$$

$$y''(t) = R\varphi''(t) \cos \varphi(t) - R(\varphi'(t))^2 \sin \varphi(t)$$

Assume there is no angular acceleration $\alpha(t) =: \varphi''(t) = 0$. That means that the angular velocity $\omega(t) =: \varphi'(t) = \text{const.} = \omega$. The angles varies linearly in time $\varphi(t) = \omega t + \phi$.

Then the velocity & the acceleration have the components on the 2 axis

$$x'(t) = -R\omega \sin(\omega t + \phi)$$

$$y'(t) = R\omega \cos(\omega t + \phi)$$

$$x''(t) = -R\omega^2 \cos(\omega t + \phi)$$

$$y''(t) = -R\omega^2 \sin(\omega t + \phi)$$

The velocity's modulus is

$$|\vec{v}(t)| = R\omega$$

While the acceleration's

$$|\vec{a}(t)| = R\omega^2$$

The way i pictured the problem, the motion is anticlockwise. Define the **centripetal acceleration in an anticlockwise uniform circular movement** as the vector which has the components

$$\vec{a}_{cp}(t) = -R\omega^2 \cos(\omega t + \phi) \vec{i} - R\omega^2 \sin(\omega t + \phi) \vec{j}$$

The velocity vector has the structure

$$\vec{v}(t) = -R\omega \sin(\omega t + \phi) \vec{i} + R\omega \cos(\omega t + \phi) \vec{j}$$

The position vector has the structure

$$\vec{r}(t) = R \cos(\omega t + \phi) \vec{i} + R \sin(\omega t + \phi) \vec{j}$$

Using these 3 last relations, u can easily prove that

$\vec{r}(t) \uparrow \downarrow \vec{a}_{cp}(t)$ and $\vec{v}(t) \perp \vec{a}_{cp}(t)$, which are typical for a circular movement. Due to the antiparallelism, the centripetal acceleration always points toward the center of the circle, hence the name.

Using the relations for the modulus of the velocity vector & the centripetal acceleration vector, u can easily prove that

$$|\vec{a}_{cp}| = \frac{|\vec{v}|^2}{R}$$

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