

This is not a very famous effect but it is really amazing. Consider a pendulum which consists of a massless rod of length r and a point of mass m ; the system is in the standard gravitational field \mathbf{g} . So the point m moves on the sphere of radius r . It is clear, the equilibrium when the point rests in the North Pole of the sphere is unstable.

Introduce a Cartesian inertial frame $OXYZ$ with origin in the point of suspension and the axis OZ is vertical such that $\mathbf{g} = -g\mathbf{e}_z$. Now let us switch on a Lorentz force $\mathbf{F} = \mathbf{B} \times \mathbf{v}$ which acts on m . The vector $\mathbf{B} = B\mathbf{e}_z$ is constant.

Theorem. Assume that B is sufficiently big:

$$\frac{B^2}{8m} > \frac{mg}{2r}$$

then the North Pole equilibrium is stable.

FIGURE 1. ...

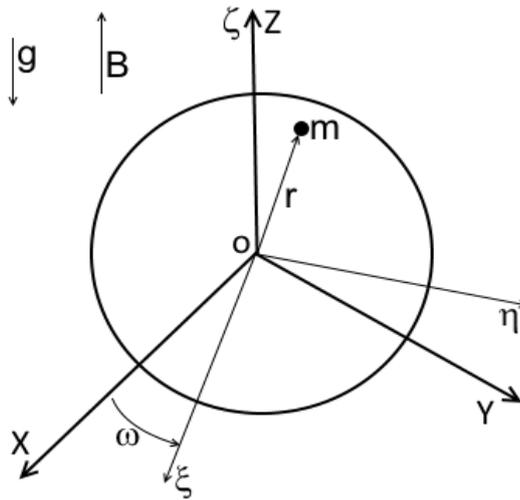


FIGURE 2. ...

The equation of particle's motion is

$$m\mathbf{a} = \mathbf{N} + m\mathbf{g} + \mathbf{B} \times \mathbf{v}, \quad (*)$$

here $\mathbf{N} = N\mathbf{r}/r$ – is a reaction of the rod.

Introduce a moving frame $O\xi\eta\zeta$, such that axis $O\zeta$ coincides with Oz . The frame $O\xi\eta\zeta$ rotates about the axis $O\zeta$ with angular velocity $\boldsymbol{\omega} = \omega\mathbf{e}_z$. The constant ω to be determined.

Let $\mathbf{a}_e, \mathbf{v}_e$ stand for the acceleration and the velocity of the moving frame; \mathbf{a}_c is the Coriolis acceleration; $\mathbf{v}_r, \mathbf{a}_r$ are the velocity and acceleration of the particle relatively the frame $O\xi\eta\zeta$. We have

$$\mathbf{a} = \mathbf{a}_r + \mathbf{a}_c + \mathbf{a}_e, \quad \mathbf{v} = \mathbf{v}_r + \mathbf{v}_e,$$

2

and

$$\mathbf{a}_e = \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r}), \quad \mathbf{a}_c = 2\boldsymbol{\omega} \times \mathbf{v}_r, \quad \mathbf{v}_e = \boldsymbol{\omega} \times \mathbf{r}.$$

Substitute these to the equation of motion

$$m\mathbf{a}_r = \mathbf{N} + m\mathbf{g} + \mathbf{B} \times (\mathbf{v}_r + \boldsymbol{\omega} \times \mathbf{r}) - m(\boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r}) + 2\boldsymbol{\omega} \times \mathbf{v}_r).$$

To get rid from the gyroscopic terms in the last formula choose

$$\boldsymbol{\omega} = \frac{1}{2m}\mathbf{B}.$$

It follows that

$$m\mathbf{a}_r = \mathbf{N} + m\mathbf{g} + \frac{1}{4m}\mathbf{B} \times (\mathbf{B} \times \mathbf{r}).$$

The last equation possesses the energy integral

$$\frac{1}{2}m|\mathbf{v}_r|^2 + V(\mathbf{r}), \quad V(\mathbf{r}) = -m(\mathbf{g}, \mathbf{r}) + \frac{1}{8m}|\mathbf{B} \times \mathbf{r}|^2.$$

This is checked by direct calculation. In the vicinity of the North Pole we have $\mathbf{r} = \xi\mathbf{e}_\xi + \eta\mathbf{e}_\eta + \sqrt{r^2 - \xi^2 - \eta^2}\mathbf{e}_\zeta$, and

$$V = \left(\frac{B^2}{8m} - \frac{mg}{2r}\right)\rho^2 + o(\rho^2),$$

as $\rho^2 = \xi^2 + \eta^2 \rightarrow 0$. QED