



Рис. 1. ...

Let S be the centre of hoop and let φ be the angle between the horizontal plane and hoop's plane. The letter A stands for the contact point of the hoop and the pillar.

Introduce a moving Cartesian frame $Sxyz$. The axis Sx remains horizontal (we are looking for such a motion) and belongs to the hoop's plane; the axis Sy paths through the point

$$A, \quad \mathbf{SA} = R\mathbf{e}_y.$$

Correspondingly, the axis Sz is perpendicular to the plane of hoop.

Then $\mathbf{e} = \cos \varphi \mathbf{e}_z + \sin \varphi \mathbf{e}_y$ is the unit vertical vector, $\mathbf{g} = -g\mathbf{e}$. Let \mathbf{T} stand for the reaction force of the pillar.

The equations of hoop's motion are

$$J_S \boldsymbol{\varepsilon} + [\boldsymbol{\omega}, J_S \boldsymbol{\omega}] = [\mathbf{SA}, \mathbf{T}], \quad m\mathbf{a}_S = m\mathbf{g} + \mathbf{T}.$$

Here m is the mass of hoop; $J_S = mR^2 \text{diag}(1/2, 1/2, 1)$ is the hoop's inertia tensor with respect to the point S ; $\boldsymbol{\varepsilon}, \boldsymbol{\omega}$ are hoop's angular acceleration and angular velocity respectively; \mathbf{a}_S is the acceleration of the hoop's centre.

From these two equation one obtains

$$J_S \boldsymbol{\varepsilon} + [\boldsymbol{\omega}, J_S \boldsymbol{\omega}] = m[\mathbf{SA}, \mathbf{a}_S - \mathbf{g}]. \quad (*)$$

The angular velocity of the frame $Sxyz$ is expressed as follows $\boldsymbol{\omega}_e = \Omega \mathbf{e}$; and $\boldsymbol{\omega}_r = \nu \mathbf{e}_z$ is the angular velocity of the hoop in its motion relatively to the frame $Sxyz$.

We are looking for the motion such that the values ν, Ω, φ are constants.

Thus the absolute angular velocity of the hoop is given by the formula $\boldsymbol{\omega} = \boldsymbol{\omega}_e + \boldsymbol{\omega}_r$. The absolute angular acceleration is $\boldsymbol{\varepsilon} = [\boldsymbol{\omega}_e, \boldsymbol{\omega}_r] = \Omega \nu \sin \varphi \mathbf{e}_x$.

Introduce a horizontal vector

$$\mathbf{PS} = (R \cos \varphi - r)(-\cos \varphi \mathbf{e}_y + \sin \varphi \mathbf{e}_z), \quad R \cos \varphi - r > 0.$$

Then the velocity of the point S is $\mathbf{v}_S = [\boldsymbol{\omega}_e, \mathbf{PS}]$, and

$$\mathbf{a}_S = [\boldsymbol{\omega}_e, \frac{d}{dt} \mathbf{PS}] = [\boldsymbol{\omega}_e, [\boldsymbol{\omega}_e, \mathbf{PS}]].$$

Substituting all these formulas to (*), we yield

$$2R\Omega\nu\sin\varphi + \Omega^2\sin\varphi(3R\cos\varphi - 2r) - 2g\cos\varphi = 0. \quad (**)$$

Non-slip condition $\mathbf{v}_S = [\boldsymbol{\omega}, \mathbf{A}\mathbf{S}] = [\boldsymbol{\omega}_e, \mathbf{P}\mathbf{S}]$ implies $\nu = -r\Omega/R$. Substitute this formula to (**) and find

$$\Omega^2 = \frac{2g\cos\varphi}{\sin\varphi(3R\cos\varphi - 4r)}.$$

Consequently, one has $3R\cos\varphi - 4r > 0$, particularly this gives $R > 4r/3$.

The normal component to pillar's surface of the vector \mathbf{T} is

$$T_n = m\Omega^2(R\cos\varphi - r).$$

The tangent component is $T_\tau = mg$. So that we have $T_\tau < \mu T_n$, here μ is the friction coefficient.