

Ricci Tensor - the amount by which the volume element of a geodesic ball in a curved Riemannian manifold deviates from the standard ball in Euclidean space.

$$T_{ss} \approx T_{\text{expl}}$$

$T_{\text{hydr}}$  hydrostatic time-scale

$$T_{\text{hydr}} \approx \left( \frac{R}{\frac{GM}{R^2}} \right)^{\frac{1}{2}} \approx \left( \frac{R^3}{GM} \right)^{\frac{1}{2}}$$

## 2.5 Non-Spherical Case

$$e \frac{dv}{dt} = -\nabla P - e \nabla \Phi (r(v, t)) \quad \text{Eulerian} \quad -\frac{1}{r^2} = -\nabla$$

$v$  is the velocity vector, and substantial time derivative on the left is defined by the operator:

$$\frac{d}{dt} = \frac{\partial}{\partial t} + v \cdot \nabla$$

The general form of ~~the~~ has already been shown to be the continuity equation of hydrodynamics:

$$\frac{\partial e}{\partial t} = -\nabla \cdot (ev)$$

Poisson equation ~~the~~ - gravitational potential ( $\Phi$ ) is connected with an arbitrary distribution of the density:

$$\nabla^2 \Phi = 4\pi G e$$

## 2.6 Hydrostatic Equilibrium in General Relativity

Einstein field equations:

$$R_{ik} - \frac{1}{2} g_{ik} R = \frac{K}{c^2} T_{ik}, \quad K = \frac{8\pi G}{c^2}$$

$R_{ik}$  = Ricci Tensor

$g_{ik}$  = metric tensor

$R$  = Riemann Curvature

$T_{ik}$  = Energy-momentum tensor

$T_{00} = e c^2, T_{11} = T_{22} = T_{33} = P$  ( $e$  includes the energy density,  $P$  = pressure)

Line element  $ds$  (distance between two neighbouring events) is given in spherical coordinates  $(r, \vartheta, \varphi)$  (time-independent)

$$ds^2 = e^{\nu} c^2 dt^2 - e^{\lambda} dr^2 - r^2 (d\vartheta^2 + \sin^2 \vartheta d\varphi^2)$$

with  $v = v(r)$ ,  $\lambda = \lambda(r)$ . Expressions  $T_{ik}$  and  $ds$ ; the field (2.24) equations can be reduced to 3 ordinary differential equations:

$$\frac{K\rho}{c^2} = e^{-\lambda} \left( \frac{v'}{r} + \frac{1}{r^2} \right) - \frac{1}{r^2},$$

$$\frac{K\rho}{c^2} = \frac{1}{2} e^{-\lambda} \left( v'' + \frac{1}{2} v'^2 + \frac{v' - \lambda'}{r} - \frac{v\lambda'}{2} \right),$$

$$K\rho = e^{-\lambda} \left( \frac{\lambda'}{r} - \frac{1}{r^2} \right) + \frac{1}{r^2}$$

Primes denote derivatives with respect to  $r$

Multiply  $4\pi r^2$  with  $\uparrow = Km = 4\pi r(1 - e^{-\lambda})$

Here  $m$  denotes the "gravitational mass" inside  $r$ :

$$m = \int_0^r 4\pi r^2 \rho dr$$

$$\rho = \rho_0 + \frac{U}{c^2}$$

$U$  = whole energy density  
 $\rho_0$  = rest-mass density

The changed metric would give the spherical volume element as  $e^{\frac{\lambda}{2}} \cdot 4\pi r^2 dr$  instead of the usual form  $4\pi r^2 dr$

Differentiation of  $\frac{K\rho}{c^2} = e^{-\lambda} \left( \frac{v'}{r} + \frac{1}{r^2} \right) - \frac{1}{r^2}$  (with respect to  $r$ ) gives  $P' = P'(\lambda, \lambda', v', v'', r)$  where  $\lambda', v', v''$  are arrives at the Tolman - Oppenheimer - Volkoff (TOV) equations for hydrostatic equilibrium in general relativity:

$$\frac{dP}{dr} = - \frac{Gm}{r^2} e \left( 1 + \frac{P}{\rho c^2} \right) \left( 1 + \frac{4\pi r^3 P}{mc^2} \right) \left( 1 - \frac{2Gm}{rc^2} \right)^{-1} \quad (2.4)$$

How does one come to this equation?

$$ds^2 = e^v c^2 dt^2 - e^\lambda dr^2 - r^2 (d\vartheta^2 + \sin^2 \vartheta d\varphi^2)$$

$$T_{ik} = T_{00} = \rho c^2, T_{11} = T_{22} = T_{33} = P$$

$$R = \frac{K}{c^2} T_{ik}$$

$$v = v(r), \lambda = \lambda(r) \quad \frac{K\rho}{c^2} = e^v e^{-\lambda} \left( \frac{v'}{r} + \frac{1}{r^2} \right) - \frac{1}{r^2} \rightarrow \text{with respect to } r$$

$$\frac{K\rho}{c^2} = e^v - \frac{dt^2}{r^2} + \frac{e^{-\lambda}}{r^2} = - \frac{1}{r^2} \left( \frac{d\vartheta}{r} + \frac{\sin \vartheta dp}{r^2} \right)$$

$$\frac{K\rho}{c^2} = e^{-\lambda} \left( \frac{v'}{r} + \frac{1}{r^2} \right) - \frac{1}{r^2}$$