



Figure 6-3 Construction for finding far field of dipoles 2 and 4 of square loop.

The far field of the individual dipole was developed in Chap. 5, being given in Table 5-1. In developing the dipole formula, the dipole was in the  $z$  direction, whereas in the present case it is in the  $x$  direction (see Figs. 6-2 and 6-3). The angle  $\theta$  in the dipole formula is measured from the dipole axis and is  $90^\circ$  in the present case. The angle  $\theta$  in (5) is a different angle with respect to the dipole, being as shown in Figs. 6-2 and 6-3. Thus, we have for the far field  $E_{\phi 0}$  of the individual dipole

$$E_{\phi 0} = \frac{j60\pi[I]L}{r\lambda} \quad (6)$$

where  $[I]$  is the retarded current on the dipole and  $r$  is the distance from the dipole. Substituting (6) in (5) then gives

$$E_{\phi} = \frac{60\pi[I]Ld_r \sin \theta}{r\lambda} \quad (7)$$

However, the length  $L$  of the short dipole is the same as  $d$ , that is,  $L = d$ . Noting also that  $d_r = 2\pi d/\lambda$  and that the area  $A$  of the loop is  $d^2$ , (7) becomes

$$E_{\phi} = \frac{120\pi^2[I] \sin \theta}{r} \frac{A}{\lambda^2} \quad (8)$$

This is the instantaneous value of the  $E_{\phi}$  component of the far field of a small loop of area  $A$ . The peak value of the field is obtained by replacing  $[I]$  by  $I_0$ , where  $I_0$  is the peak current in time on the loop. The other component of the far field of the loop is  $H_{\theta}$ , which is obtained from (8) by dividing by the intrinsic impedance of the medium, in this case, free space. Thus,

$$H_{\theta} = \frac{E_{\phi}}{120\pi} = \frac{\pi[I] \sin \theta}{r} \frac{A}{\lambda^2} \quad (9)$$