

Expect a metric of the form:

$$ds^2 = g_{tt}dt^2 + g_{rr}dr^2 \quad (1)$$

## 0.1 $\mu = r$ Component

The Lorentz equations for the  $\mu = r$  component is:

$$\frac{dp_r}{d\tau} + \Gamma_{\alpha\beta}^r p_\alpha \frac{dx^\beta}{d\tau} = -q F^{r\nu} g_{\nu\beta} \frac{dx^\beta}{d\tau}$$

Thus:

$$\begin{aligned} \frac{d}{d\tau} \frac{dr}{d\tau} + \Gamma_{tt}^r \frac{dt}{d\tau} \frac{dt}{d\tau} + \Gamma_{tr}^r \frac{dt}{d\tau} \frac{dr}{d\tau} + \Gamma_{rt}^r \frac{dr}{d\tau} \frac{dt}{d\tau} + \Gamma_{rr}^r \frac{dr}{d\tau} \frac{dr}{d\tau} \\ = -\frac{q}{m} F^{r\nu} \left( g_{\nu t} \frac{dt}{d\tau} + g_{\nu r} \frac{dr}{d\tau} \right) \end{aligned}$$

using the metric 1:

$$\begin{aligned} \frac{d}{d\tau} \frac{dr}{d\tau} + \Gamma_{tt}^r \frac{dt}{d\tau} \frac{dt}{d\tau} + \Gamma_{tr}^r \frac{dt}{d\tau} \frac{dr}{d\tau} + \Gamma_{rt}^r \frac{dr}{d\tau} \frac{dt}{d\tau} + \Gamma_{rr}^r \frac{dr}{d\tau} \frac{dr}{d\tau} \\ = -\frac{q}{m} F^{rt} \left( g_{tt} \frac{dt}{d\tau} \right) - q F^{rr} \left( g_{rr} \frac{dr}{d\tau} \right) \end{aligned}$$

Since,  $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$ , then  $F_{rt} = \partial_r A_t - \partial_t A_r = \frac{d\psi}{dr}$  and  $F_{rr} = \partial_r A_r - \partial_r A_r = 0$ , so:

$$\begin{aligned} \frac{d}{d\tau} \frac{dr}{d\tau} + \Gamma_{tt}^r \frac{dt}{d\tau} \frac{dt}{d\tau} + \Gamma_{tr}^r \frac{dt}{d\tau} \frac{dr}{d\tau} + \Gamma_{rt}^r \frac{dr}{d\tau} \frac{dt}{d\tau} + \Gamma_{rr}^r \frac{dr}{d\tau} \frac{dr}{d\tau} \\ = -\frac{q}{m} g^{tt} g^{rr} \frac{d\psi}{dr} \left( g_{tt} \frac{dt}{d\tau} \right) \quad (2) \end{aligned}$$

Now, the Christoffel symbols we require are  $\Gamma_{tt}^r, \Gamma_{tr}^r (= \Gamma_{rt}^r), \Gamma_{rr}^r$  and are given by:

$$\Gamma_{tt}^r = \frac{1}{2} g^{r\beta} (\partial_t g_{\beta t} + \partial_t g_{\beta t} - \partial_\beta g_{tt}) = \frac{1}{2} g^{rr} (\partial_t g_{rt} + \partial_t g_{tr} - \partial_r g_{tt}) = \frac{1}{2} g^{rr} (-\partial_r g_{tt})$$

$$\Gamma_{tr}^r = \frac{1}{2} g^{r\beta} (\partial_r g_{\beta t} + \partial_t g_{\beta r} - \partial_\beta g_{tr}) = \frac{1}{2} g^{rr} (\partial_r g_{rt} + \partial_t g_{rr} - \partial_r g_{tr}) = 0$$

$$\Gamma_{rr}^r = \frac{1}{2} g^{r\beta} (\partial_r g_{\beta r} + \partial_r g_{\beta r} - \partial_\beta g_{rr}) = \frac{1}{2} g^{rr} (\partial_r g_{rr} + \partial_r g_{rr} - \partial_r g_{rr}) = \frac{1}{2} g^{rr} (\partial_r g_{rr})$$

Substituting these Christoffel symbols into 2:

$$\frac{d}{d\tau} \frac{dr}{d\tau} + \frac{1}{2} g^{rr} (-\partial_r g_{tt}) \frac{dt}{d\tau} \frac{dt}{d\tau} + \frac{1}{2} g^{rr} (\partial_r g_{rr}) \left( \frac{dr}{d\tau} \right)^2 = -\frac{q}{m} g^{tt} g^{rr} \frac{d\psi}{dr} \left( g_{tt} \frac{dt}{d\tau} \right)$$

$$\Rightarrow \frac{d}{d\tau} \frac{dr}{d\tau} - \frac{1}{2} g^{rr} \partial_r g_{tt} \left( \frac{dt}{d\tau} \right)^2 + \frac{1}{2} g^{rr} \partial_r g_{rr} \left( \frac{dr}{d\tau} \right)^2 = -\frac{q}{m} g^{rr} \frac{d\psi}{dr} \frac{dt}{d\tau} \quad (3)$$

The proper time,  $\tau$ , is given by:

$$d\tau^2 = g_{tt} dt^2 + g_{rr} dr^2$$

Rearranging:

$$\begin{aligned} 1 &= g_{tt} \left( \frac{dt}{d\tau} \right)^2 + g_{rr} \left( \frac{dr}{d\tau} \right)^2 \\ \Rightarrow \left( \frac{dt}{d\tau} \right)^2 &= \frac{1}{g_{tt}} - \frac{g_{rr}}{g_{tt}} \left( \frac{dr}{d\tau} \right)^2 \end{aligned} \quad (4)$$

Substitute 4 into 3:

$$\begin{aligned} \frac{d}{d\tau} \frac{dr}{d\tau} - \frac{1}{2} g^{rr} \partial_r g_{tt} \left( \frac{1}{g_{tt}} - \frac{g_{rr}}{g_{tt}} \left( \frac{dr}{d\tau} \right)^2 \right) + \frac{1}{2} g^{rr} \partial_r g_{rr} \left( \frac{dr}{d\tau} \right)^2 &= -\frac{q}{m} g^{rr} \frac{d\psi}{dr} \frac{dt}{d\tau} \\ \Rightarrow \frac{d}{d\tau} \frac{dr}{d\tau} - \frac{g^{rr} \partial_r g_{tt}}{2g_{tt}} + \frac{g^{rr} g_{rr} \partial_r g_{tt}}{2g_{tt}} \left( \frac{dr}{d\tau} \right)^2 + \frac{g^{rr} \partial_r g_{rr}}{2} \left( \frac{dr}{d\tau} \right)^2 &= -\frac{q}{m} g^{rr} \frac{d\psi}{dr} \frac{dt}{d\tau} \\ \Rightarrow \frac{d}{d\tau} \frac{dr}{d\tau} + \left( \frac{\partial_r g_{tt}}{2g_{tt}} + \frac{\partial_r g_{rr}}{2g_{tt}} \right) \left( \frac{dr}{d\tau} \right)^2 - \frac{\partial_r g_{tt}}{2g_{rr} g_{tt}} &= -\frac{q}{m} g^{rr} \frac{d\psi}{dr} \frac{dt}{d\tau} \end{aligned}$$