

1. The Line Element

In cylindrical coordinates t, z, r, ϕ ,

$$ds^2 = dt^2 (L^2 - 1) + 2 r L d\phi dt + r^2 d\phi^2 + dz^2 + dr^2$$

where $0 < L < 1$ is a constant.

The proper acceleration of the vector

$$U_\mu = \sqrt{1 - L^2} \vec{\partial}_t - r L \vec{\partial}_\phi / \sqrt{1 - L^2}$$

is $\dot{U}_\nu = (\nabla_\mu U_\nu) U^\mu$, the covariant derivative of U_ν projected in the direction U^μ . The covector $g^{\mu\nu} U_\nu$ has only a time (first) component so the only terms of $D_{\mu\nu} = \nabla_\mu U_\nu$ which can be selected by the contraction of $D_{\mu\nu}$ with U^ν are in the first column, $D_{1,k}$.

Using $\nabla_\mu U_\nu = \partial_\mu U_\nu - \Gamma_{\mu\nu}^\alpha U_\alpha$ there are eight instances where $\Gamma_{\mu\nu}^\alpha U_\alpha$ is non-zero listed below in the 'Trace output section' in the pdf.

The first two lines are equivalent to

$$\begin{aligned} D_{12} &= -\Gamma_{12}^1 U_1 - \Gamma_{12}^4 U_4 \\ &= \frac{L (L^3 - L)}{2 r \sqrt{1 - L^2}} - \frac{(L - 1) L^2 (L + 1)}{2 r \sqrt{1 - L^2}} \\ &= 0 \end{aligned}$$

and since there are no other candidates in the first column of $D_{\mu\nu}$ this is sufficient to eliminate any proper acceleration.

2. Trace output

The amended script reports this as it loops through μ, ν, α . These are the first three digits.

$$\begin{aligned} 121 \quad dU = 0 \quad U(1) &= -\frac{(L - 1) (L + 1)}{\sqrt{1 - L^2}} \quad mcs = \frac{L^2}{2r} >> -\frac{(L - 1) L^2 (L + 1)}{2 r \sqrt{1 - L^2}} \\ 124 \quad dU = 0 \quad U(4) &= -\frac{r L}{\sqrt{1 - L^2}} \quad mcs = -\frac{L^3 - L}{2 r^2} >> \frac{L (L^3 - L)}{2 r \sqrt{1 - L^2}} \\ 211 \quad dU = 0 \quad U(1) &= -\frac{(L - 1) (L + 1)}{\sqrt{1 - L^2}} \quad mcs = \frac{L^2}{2r} >> -\frac{(L - 1) L^2 (L + 1)}{2 r \sqrt{1 - L^2}} \\ 214 \quad dU = 0 \quad U(4) &= -\frac{r L}{\sqrt{1 - L^2}} \quad mcs = -\frac{L^3 - L}{2 r^2} >> \frac{L (L^3 - L)}{2 r \sqrt{1 - L^2}} \\ 241 \quad dU = 0 \quad U(1) &= -\frac{(L - 1) (L + 1)}{\sqrt{1 - L^2}} \quad mcs = \frac{L}{2} >> -\frac{(L - 1) L (L + 1)}{2 \sqrt{1 - L^2}} \\ 244 \quad dU = 0 \quad U(4) &= -\frac{r L}{\sqrt{1 - L^2}} \quad mcs = -\frac{L^2 - 2}{2 r} >> \frac{L (L^2 - 2)}{2 \sqrt{1 - L^2}} \\ 421 \quad dU = -\frac{L}{\sqrt{1 - L^2}} \quad U(1) &= -\frac{(L - 1) (L + 1)}{\sqrt{1 - L^2}} \quad mcs = \frac{L}{2} >> -\frac{(L - 1) L (L + 1)}{2 \sqrt{1 - L^2}} \\ 424 \quad dU = -\frac{L}{\sqrt{1 - L^2}} \quad U(4) &= -\frac{r L}{\sqrt{1 - L^2}} \quad mcs = -\frac{L^2 - 2}{2 r} >> \frac{L (L^2 - 2)}{2 \sqrt{1 - L^2}} \end{aligned}$$

References

- [1] Leslie E. Ballentine
Quantum Mechanics - A Modern Development
World Scientific (1998)